



MICROCOFY

CHART

MRC Technical Summary Report #2901

ON THE EXISTENCE OF PERIODIC SOLUTIONS FOR A CLASS OF SYMMETRIC HAMILTONIAN SYSTEMS

Paul H. Rabinowitz



Mathematics Research Center University of Wisconsin—Madison 610 Walnut Street Madison, Wisconsin 53705

January 1986

(Received December 3, 1985)



Approved for public release Distribution unlimited

Sponsored by

U. S. Army Research Office P. O. Box 12211 Research Triangle Park North Carolina 27709 National Science Foundation Washington, DC 20550

86 5 20 152

UNIVERSITY OF WISCONSIN-MADISON MATHEMATICS RESEARCH CENTER

ON THE EXISTENCE OF PERIODIC SOLUTIONS FOR A CLASS OF SYMMETRIC HAMILTONIAN SYSTEMS

Paul H. Rabinowitz

Technical Summary Report #2901 January 1986

ABSTRACT

The main result in this paper is:

Theorem: If $H \in C^1(\mathbb{R}^{2n}, \mathbb{R})$ and satisfies

 (H_1) $H^{-1}(1)$ bounds a starshaped neighborhood of 0 in \mathbb{R}^{2n} ,

 (H_2) z • $H_z \neq 0$ for all z $\in H^{-1}(1)$,

(H₃) H(p,q) = H(-p,q) for all $p,q \in \mathbb{R}^n$, then there is a T > 0 such that the Hamiltonian system

(HS)
$$\dot{z} = H_z(z), = \begin{pmatrix} 0 & -id \\ id & 0 \end{pmatrix}$$

possesses a T periodic solution $(p(t),q(t)) \in H^{-1}(1)$ with p odd about 0 and T/2 and q even about 0 and T/2.

The proof involves a new existence mechanism which should be useful in other situations.

AMS (MOS) Subject Classifications: 34C25, 58E05, 58F22, 70H05, 70H25, 70K99

Key Words: periodic solution, Hamiltonian system, minimax methods

Work Unit Number 1 (Applied Analysis)

Mathematics Department and Mathematics Research Center, University of Wisconsin-Madison, Madison, WI 53705.

This research was sponsored in part by the United States Army under Contract No. DAAG29-80-C-0041 and by the National Science Foundation under Grant No. MCS-8110556. Reproduction in whole or in part for any purpose of the United States Government is permitted.

SIGNIFICANCE AND EXPLANATION

Hamiltonian systems are used to model the motion of discrete mechanical systems. This paper establishes the existence of periodic solutions for a class of such systems. The method developed to prove existence should be useful for other such problems.

Accesion For				
	CRA&I	A]	
DTIC	TAB	F		
Ur.announced				
Justification				
By				
Availability Codes				
Dist	Avail and or Special			
4-1				
ו־ח				



The responsibility for the wording and views expressed in this descriptive summary lies with MRC, and not with the author of this report.

ON THE EXISTENCE OF PERIODIC SOLUTIONS FOR A CLASS OF SYMMETRIC HAMILTONIAN SYSTEMS Paul H. Rabinowitz

1. Introduction

Consider the Hamiltonian system of ordinary differential equations:

(HS)
$$\dot{z} = JH_{z}(z), \quad J = \begin{pmatrix} 0 & -id \\ id & 0 \end{pmatrix}$$

Here $H: \mathbb{R}^{2n} + \mathbb{R}$, z = (p,q) with $p,q \in \mathbb{R}^n$, and id denotes the $n \times n$ identity matrix. Several papers have investigated what conditions on H lead to the existence of periodic solutions of (HS) having prescribed energy, i.e. H(z) is a given constant. See e.g. [1-11]. (Other studies such as [12] treat the multiplicity of periodic solutions of (HS) of prescribed energy.) In particular, it was shown in [4] that Theorem 1.1: If $H \in C^1(\mathbb{R}^{2n}, \mathbb{R})$ and satisfies

 (H_1) $H^{-1}(1)$ is the boundary of a starshaped neighborhood of 0 in \mathbb{R}^{2n} ,

 (H_2) z • $H_z(z) \neq 0$ on $H^{-1}(1)$,

Madison, WI 53705.

then (HS) possesses a periodic solution on $H^{-1}(1)$.

In Theorem 1.1, "starshaped" means $H^{-1}(1)$ is homeomorphic to S^{2n-1} by a radial projection map.

Our goal in this paper is to show that if H satisfies an additional symmetry condition, (HS) possesses a periodic solution having additional properties: $\frac{\text{Theorem 1.2:}}{\text{Theorem 1.2:}} \quad \text{If } \text{H } \in \text{C}^1(\mathbb{R}^{2n},\mathbb{R}) \quad \text{and satisfies } (\text{H}_1)^-(\text{H}_2) \quad \text{and} \quad \text{(H}_3) \quad \text{H}(\text{p,q}) = \text{H}(-\text{p,q}) \quad \text{for all} \quad \text{p,q} \in \mathbb{R}^n \quad ,$ then there exists a T > 0 and a T periodic solution $(\text{p(t),q(t)}) \quad \text{of (HS) on}$

 $H^{-1}(1)$ such that p is odd and q is even about t=0 and $\frac{T}{2}$.

Mathematics Department and Mathematics Research Center, University of Wisconsin-Madison,

This research was sponsored in part by the United States Army under Contract No. DAAG29-80-C-0041 and by the National Science Foundation under Grant No. MCS-8110556. Reproduction in whole or in part for any purpose of the United States Government is permitted.

Periodic solutions of this type were studied by SEIFERT [1], RUIZ [2], WEINSTEIN [3], GLUCK-ZILLER [7], HAYASHI [8], and BENCI [9] for C^2 or smoother H's satisfying (H_3) and having the form H(p,q) = K(p,q) + V(q) with K and V suitably restricted. Different symmetries for (HS) have been treated by van GROESEN [10] and GIRARDI [11]. In [10] it was shown that if $H \in C^2$, satisfies (H_2) , $H^{-1}(1)$ bounds a convex region and H(p,q) = H(-p,q) = H(p,-q) for all $p,q \in \mathbb{R}^n$, then the conclusions of Theorem 1.2 hold. The convexity assumption on $H^{-1}(1)$ plays a strong role in the existence argument here. In [11] on the other hand, $(H_1)^-(H_2)$ and H(z) = H(-z) are assumed and it is proved that there is a $\tau > 0$ such that (HS) possesses a 2τ periodic solution on $H^{-1}(1)$ for which $z(t + \tau) = -z(\tau)$ for all $t \in \mathbb{R}$.

Both [4] and [11] rely on minimax arguments and topological index theories to exploit an S^1 symmetry associated with a variational formulation of (HS). Topological index theories are often useful in obtaining multiple critical points of a symmetric functional and indeed such multiplicity results were the main goal of [10-11] and enabled them to obtain analogues of a theorem of EKELAND and LASRY [12]. For the problem treated here however, we will work directly in the space of T periodic functions p,q for which p is odd about t=0 and T/2 and q is even about 0 and T/2. The symmetries used earlier in [4,10,11] are not longer present if one works in this space and therefore another existence mechanism to treat (HS) is required. Indeed developing such a new mechanism is one of our main goals here. In a future note, we will show how this method can also be applied to treat the sort of situation studied in [1-3,5, etc.].

In §2, the solution of (HS) will be reduced to finding a critical point of an associated variational problem. Existence of a critical point when H \in C² is carried out in §3. Lastly §4 contains the C¹ case as well as the proof of a crucial intersection theorem used in §3. For some of the technicalities of §2-4, we have benefited from unpublished work of V. BENCI and the author.

§2. Formulation of the Variational Problem

In this section the solution of (HS) will be reduced to finding critical points of a variational problem. For technical convenience we assume for now that $H \in C^2(\mathbb{R}^{2n}, \mathbb{R})$. The more general case of $H \in C^1(\mathbb{R}^{2n}, \mathbb{R})$ will be treated in §4.

The variational formulation of (HS) will take place in an "L²" type of setting and therefore the behavior of H outside of H⁻¹(1) is important. Thus as in [4] we define a new Hamiltonian $\overline{H}(z)$ which coincides with H on H⁻¹(1) and grows at a controlled rate as $|z| + \infty$. Since H⁻¹(1) bounds a starshaped region, for all $z \in \mathbb{R}^{2n} \setminus \{0\}$, there is a unique $\alpha(z) > 0$ and $w(z) \in H^{-1}(1)$ such that $z = \alpha(z)w(z)$. In fact, w depends only on $z|z|^{-1}$ and $\alpha(z) = |z||w(z)|^{-1}$. Define $\overline{H}(0) = 0$ and for $z \neq 0$, $\overline{H}(z) = \alpha(z)^2$. It is easy to check that $\overline{H} \in \mathbb{C}^2(\mathbb{R}^{2n} \setminus \{0\}, \mathbb{R})$, \overline{H}_{ZZ} is uniformly bounded, and \overline{H} is homogeneous of degree two. Moreover $\overline{H}^{-1}(1) = H^{-1}(1)$ and $\overline{H}(p,q)$ is even in p. For future reference, note the following estimates for \overline{H} . Let

Then by the homogeneity of \overline{H} , for all $z \in \mathbb{R}^{2n}$,

(2.2)
$$m|z|^2 \le \overline{H}(z) \le M|x|^2$$
.

Lemma 2.3: Suppose there exists $\lambda > 0$ and a 2π periodic solution $\zeta(\tau) = (\phi(\tau), \psi(\tau))$

$$\dot{\xi} = \lambda J \vec{H}_{\mu}(\zeta)$$

with $\zeta(\tau) \in \widetilde{\mathbb{H}}^{-1}(1)$, ϕ odd about 0 and π and ψ even about 0 and π . Then there exists a T > 0 and a T periodic solution of (HS) of the type stated in Theorem 1.2. Proof: Set $z(t) = \zeta(r(t))$ where $r \in C^1$ is free for the moment. Then z is a solution of (HS) if

(2.5)
$$\dot{z} = \lambda \vec{H}_{z}(\zeta(r(t)))\dot{r} = JH_{z}(\zeta(r(t)))$$
.

Since $H^{-1}(1) = \tilde{H}^{-1}(1)$ and H_z , $\tilde{H}_z \neq 0$ on this set, there is a function $\beta \in C^1(H^{-1}(1), \mathbb{R}\setminus\{0\}) \text{ such that } \tilde{H}_Z(z) = \beta(z)\tilde{H}_Z(z) \text{ for } z \in H^{-1}(1). \text{ Therefore 2.5 shows}$ $(2.6) \qquad \qquad \mathring{r} = \lambda^{-1}\beta(\zeta(r(t))).$

Further setting r(0) = 0 and noting that $\beta \neq 0$, we can assume $\beta > 0$ and r is a

strictly increasing function of t. Let $T = 2r^{-1}(\pi)$. Then the properties of ζ imply if Z(t) = (P(t),Q(t)), P(0) = 0 = P(T/2) and Q'(0) = 0 = Q'(T/2). Extending P as an odd function p about 0 and T/2 and Q as an even function q about 0 and T/2, it follows that the resulting function z = (p,q) is a T periodic solution of (HS) on $H^{-1}(1)$ of the desired type.

Thus Lemma 2.3 reduces the proof of Theorem 1.2 to finding $\lambda > 0$ and a 2π periodic solution ζ of (2.4). We will convert this question to that of solving a variational problem. First an appropriate function space must be introduced. Let $x = \{z = (p,q) \in \mathbb{W}^{\frac{1}{2}}, {2 \choose 3}, \mathbb{R}^{2n}) \mid p \text{ is odd about } 0 \text{ and } \pi, \text{ } q \text{ is even about } 0 \text{ and } \pi \}.$ Here $\mathbb{W}^{\frac{1}{2}}, {2 \choose 3}, \mathbb{R}^{2n}$ is the set of 2π periodic functions

$$z = \sum_{j \in \mathbb{Z}} a_j e^{ijt}$$

such that

TOOK THURSDAY THE PROPERTY OF THE PROPERTY OF

$$\sum_{j \in \mathbb{Z}} (1 + |j|) |a_j|^2 < \infty.$$

For smooth $z \in X$, let

(2.7)
$$A(z) = \int_{0}^{2\pi} p \cdot \dot{q} dt$$
.

Then

$$|A(z)| \le \text{const.} |z|^2 \frac{1}{2}, 2$$

i.e. A is a continuous quadratic form on this (dense) subspace of X. Therefore A extends continuously to all of X. This extension will still be denoted by A(z).

Let e_1, \ldots, e_{2n} denote the usual basis in \mathbb{R}^{2n} , i.e. $e_1 = (1,0,0,\ldots)$, etc. and set

 $X_0 \equiv span\{e_k \mid n+1 \leq k \leq 2n\}$,

 $x^{+} \equiv span\{(sin jt)e_{k} - (cos jt)e_{k+n} | 1 \le j < \infty, 1 \le k \le n\}$,

 $X^{-} \equiv \operatorname{span}\{(\sin jx)e_{k} + (\cos jt)e_{k+n} | 1 \leq j < \infty, 1 \leq k \leq n\}.$

These spaces are mutually orthogonal in $L^2(S^1, \mathbb{R}^{2n})$. Moreover $X = X^0 \oplus X^+ \oplus X^-$ and if $z = z^0 + z^+ + z^- \in X$.

(2.8)
$$|z|^2 = |z^0|^2 + \lambda(z^+) - \lambda(z^-)$$

defines a norm on X which is equivalent to the $W^{\frac{1}{2},2}(S^1,R^{2n})$ norm. (See e.g. [4].) Setting

(2.9)
$$\Psi(z) = \frac{1}{2\pi} \int_{0}^{2\pi} \vec{H}(z) dt$$
,

the upper bound for H in (2.2) implies Y is well defined on X.

Proposition 2.10: With $\frac{\pi}{4}$ as above,

(i)
$$Y \in C^{1,Lip}(X,R)$$
,

(ii) Y' is compact.

Proof: (i) Since $\bar{H} \in C^2(\mathbb{R}^{2n} \setminus \{0\}, \mathbb{R})$ and \bar{H}_{ZZ} is uniformly bounded, there is a constant $M_1 > 0$ such that

(2.11)
$$|\vec{H}(z+\zeta) - \vec{H}(z) - \vec{H}_z(z) \cdot \zeta| \leq M_1 |\zeta|^2$$

for all $z, \zeta \in \mathbb{R}^{2n}$. Therefore for $z, \zeta \in X$, (2.11) and the continuous embedding of X in $L^2(S^1, \mathbb{R}^{2n})$ imply that

(2.12)
$$|\Psi(z+\zeta)-\Psi(z)| = \frac{1}{2\pi} \int_{0}^{2\pi} \vec{R}_{z}(z) \cdot \zeta \, dt | \leq M_{1} |\zeta|_{L^{2}}^{2} \leq M_{2} |\zeta|^{2}$$
.

Thus (2.12) shows that $Y \in C^1(X,R)$ and

(2.13)
$$\Psi^{\dagger}(z)\zeta = \frac{1}{2\pi} \int_{0}^{2\pi} \vec{H}_{z}(z) \cdot \zeta \, dt .$$

To see that "' is Lipschitz continuous, observe that

(2.14)
$$\|\Psi'(z+w) - \Psi'(z)\|_{x} = \sup_{\zeta \in X, \, \|\zeta\| \le 1} \left| \frac{1}{2\pi} \int_{0}^{2\pi} (\bar{H}_{z}(z+w) - \bar{H}_{z}(z)) \cdot \zeta dt \right|.$$

Since \vec{H}_{zz} is uniformly bounded on $\mathbb{R}^{2n}\setminus\{0\}$, there is a constant $M_3>0$ such that $|\vec{H}_z(z+w)-\vec{H}_z(z)|\leq M_3|w|$

for all $z, w \in \mathbb{R}^{2n}$. Hence (2.14)-(2.15) imply

(2.16)
$$\|\Psi'(z+w) - \Psi'(z)\|_{X}^{2} \le M_{3} \sup_{\zeta \in X, \|\zeta\| \le 1} \frac{1}{2\pi} \int_{0}^{2\pi} |w| |\zeta| dt \le M_{4} |w|$$

for all $z, w \in X$.

(ii) Let (z_j) be a bounded sequence in X. Since X is compactly embedded in $L^r(S^1,R^{2n})$ for all $r \in [1,\infty)$, (see e.g. the argument for an analogous situation in [13]), along a subsequence, z_j converges in $L^2(S^1,R^{2n})$ to $z \in X$. Hence by (2.16), (2.17) $\|\Psi'(z_j) - \Psi'(z)\|_{V^*} \leq M_4 \|z_j - z\|_{T^2} + 0$

and Y' is compact.

Let $M \equiv \Psi^{-1}(1)$.

Proposition 2.18: (i) M is a C1, Lip manifold in X.

(ii) M bounds a starshaped neighborhood of 0 in X.

(iii) M is bounded in $L^2(S^1, \mathbb{R}^{2n})$.

Proof: For $z \in X \setminus \{0\}$, by the homogeneity of \Re ,

(2.19)
$$\Psi'(z)z = \frac{1}{2\pi} \int_{0}^{2\pi} \tilde{H}_{z}(z) \cdot z \, dt = 2\Psi(z) > 0.$$

Hence M is a manifold and (i) of Proposition 2.10 shows it is $C^{1,\operatorname{Lip}}$. Moreover M is the boundary of $\Psi^{-1}(-\infty,1)$, an open set. The homogeneity of \overline{H} shows that any ray through the origin in X meets M exactly once. Hence M bounds a starshaped region. Lastly by (2.2), if $z \in M$,

(2.20)
$$\frac{m}{2\pi} \|z\|_{L^{2}}^{2} \leq \Psi(z) = 1,$$

i.e. M is bounded in $L^2(S^1, \mathbb{R}^{2n})$.

We will find a solution of the desired type as a critical point of $A|_{M}$. This functional is said to satisfy the $(PS)^+$ condition if for any c>0, whenever (z_j) is a sequence in M such that

(2.21)
$$\lambda(z_1) + c$$

and

(2.22)
$$A_{M}^{\dagger}(z_{j}) \equiv A^{\dagger}(z_{j}) - \lambda(z_{j})\Psi^{\dagger}(z_{j}) + 0 \text{ in } X^{\dagger}$$

where

$$\lambda(z) = (\lambda^{*}(z), \Psi^{*}(z))_{X} \Psi^{*}(z) \Psi^{*}_{X},$$

then (z_j) possesses a convergent subsequence. Thus the $(PS)^+$ condition is a kind of compactness condition. It is important that $A|_{M}$ satisfy this condition in order to construct the "deformation mapping" used in §3.

Proposition 2.23: $A|_{M}$ satisfies the $(PS)^{+}$ condition.

Proof: Let $(z_j) \subset M$ and satisfy (2.21)-(2.22). Writing $z_j = z_j^+ + z_j^- + z_j^0 \in X^+ \oplus X^- \oplus X^0$, (2.22) and the homogeneity of A show

(2.24)
$$|z_{j}^{\pm}|^{2} = A'(z_{j})(\pm z_{j}^{\pm}) < |\lambda(z_{j})| \int_{0}^{2\pi} \overline{R}_{z}(z_{j}) \cdot z_{j}^{\pm} dt | + \varepsilon_{j} |z_{j}^{\pm}|$$

where $\varepsilon_j + 0$ as $j + \infty$. Since (z_j) is bounded in $L^2(S^1, \mathbb{R}^{2n})$ as is $(\overline{H}_z(z_j))$ via (2.15), by (2.24) and (2.8),

(2.25)
$$\|z_j\|^2 \le a_1(|\lambda(z_j)| + 1)$$
.

Now by the homogeneity of A and \bar{H}_{r} ,

$$(2.26) 2|a(z_{j}) - \lambda(z_{j})| = |a'(z_{j})z_{j} - \lambda(z_{j})^{\psi'}(z_{j})z_{j}| \le \varepsilon_{j}^{\sharp}z_{j}^{\sharp}.$$

Combining (2.25) and (2.26) gives

(2.27)
$$|\mathbf{A}(\mathbf{z_j}) - \lambda(\mathbf{z_j})| \leq \mathbf{a_2}(|\lambda(\mathbf{z_j})| + 1)^{1/2}$$
.

Recalling that $A(z_j) + c$, (2.27) shows $\lambda(z_j)$ is a bounded sequence and (2.25) implies (z_j) is bounded in X. Consequently (2.26) yields $\lambda(z_j) + c > 0$. Let L denote the duality map from X^* to X. Then

(2.28) $L(A^{*}(z_{j}) - \lambda(z_{j})\Psi^{*}(z_{j})) = z_{j}^{+} - z_{j}^{-} - \lambda(z_{j})L\Psi^{*}(z_{j}) + 0.$

Thus the boundedness of $(\lambda(z_j))$ and (z_j) , (ii) of Proposition 2.10, and (2.28) show z_j^+ , z_j^- , and - since x^0 is finite dimensional - z_j^0 converge along a subsequence.

§3. Existence of a Solution

In this section, the existence of a critical point of $A|_{M}$ will be established. Standard arguments then lead to a solution of (2.4) and hence (by Lemma 2.3) of (HS) of the desired type. A minimax argument will be used to get a critical point of $A|_{M}$. An important role in any minimax argument is played by the so-called deformation mapping. The following proposition lists its properties in our setting. For c > 0, let

$$A_c = \{z \in M | A(z) \le c\}$$
 and

$$K_C = \{z \in M \mid A(z) = c \text{ and } A|_H^*(z) = 0\}$$
.

Proposition 3.1: Let $c, \bar{\epsilon} > 0$. Then there exists an $\epsilon \in (0,\bar{\epsilon})$ and $\eta \in C([0,1] \times X,X)$ such that

 1° $\eta(s, \cdot)$ is a homeomorphism of X onto X for each $s \in [0, 1]$.

$$2^{\circ}$$
 $n(1,z) = z$ if $A(z) \not\in [c - \overline{\epsilon}, c + \overline{\epsilon}]$ and if $|\Psi(z) - 1| > \frac{1}{2}$.

$$3^{\circ}$$
 in(1,z) - zl < 1.

 4° n(s,M) = M for each s \in [0,1].

5° If P⁺,P⁻ denote respectively the orthogonal projectors of X onto X⁺,X⁻, then $P^{\pm}_{n(s,z)} = e^{\mp \theta(s,z)}z^{\pm} + K^{\pm}(s,z)$

where $\theta \in C([0,1] \times X, R^{+})$ and K^{\pm} is compact.

$$6^{\circ}$$
 If $K_{c} = \emptyset$, $\eta(1, A_{c+\epsilon}) \subset A_{c-\epsilon}$.

<u>Proof:</u> Most of the assertions are standard. In particular 1° and 6° as well as the precise definition of ω below can be found in [14]. (It is in proving 6° that the (PS)⁺ condition is used.) Therefore we will only verify $2^{\circ}-5^{\circ}$.

The function n satisfies an ordinary differential equation of the form

(3.2)
$$\frac{d\eta}{ds} = -\omega(\eta) \xi [A^*(\eta) - \lambda(\eta) \Psi^*(\eta)] \quad \eta(0,z) = z$$

for $z \in X$. The function $\omega \in C(X,R)$ is Lipschitz continuous and is chosen so that $0 \le \omega(z) \le 1$, the right-hand side of (3.2) in norm does not exceed 1, $\omega(z) = 0$ if $A(z) \notin \{c - \vec{\epsilon}, c + \vec{\epsilon}\}$ or if $|\Psi(z) - 1| \ge \frac{1}{2}$, and $\omega(z) \ne 0$ if $z \in M$ and A(z) is near c. This implies that (3.2) has a solution $\eta(s,z) \in C(\mathbb{R} \times X,X)$ satisfying $2^{O}-3^{O}$. The form of the right-hand side of (3.2) shows $\Psi^{1}(\eta(s,z)) \frac{d\eta}{ds} \equiv 0$ and therefore

 $\Psi(\eta(s,z)) \equiv \Psi(\eta(0,z)) = \Psi(z)$. In particular if $z \in M$, so is $\eta(s,z)$ and 4° holds. Lastly to prove 5° , note that $LA^{\dagger}(z) = z^{\dagger} - z^{-}$. Therefore $P^{\pm}\eta \cong \eta^{\pm}$ satisfies

(3.3)
$$\frac{dn^{\pm}}{ds} \pm \omega(n)n^{\pm} = \omega(n)\lambda(n)P^{\pm} L\Psi^{\dagger}(n) .$$

Treating η as being known, (3.3) shows η^{\pm} satisfies an inhomogeneous linear equation whose solution is

(3.4)
$$\eta^{\pm}(s,z) = [\exp(\mp \int_{0}^{s} \omega(\eta(r,z))dr)]z^{\pm} + \kappa^{\pm}(s,z)$$

where

$$K^{\pm}(s,z) = \int_{0}^{s} \left[\exp\left(\int_{s}^{\tau} \omega(\eta(r,z)) dr \right) \right] S^{\pm}(\eta(\tau,z)) d\tau$$

and

$$S^{\pm}(y) = \omega(y)\lambda(y)P^{\pm}L\Psi^{\dagger}(y)$$
.

Thus η is of the form asserted in 5°.

It remains only to show that K^{\pm} is compact. Note first that $S^{\pm}: X + X$ is compact. For convenience we drop the superscripts \pm for S and K. Indeed if (y_j) is bounded in X and $\Psi(y_j) \neq (\frac{1}{2}, \frac{3}{2})$ along some subsequence, then $\omega(y_j) = 0$ and $S(y_j) = 0$. Thus we can assume $\Psi(y_j) \in (\frac{1}{2}, \frac{3}{2})$ for all $j \in \mathbb{N}$. Since $\Psi(0) = 0$, this implies (y_j) is bounded away from 0. Therefore

$$\|\Psi^{\dagger}(y_{j})\|_{X^{\frac{1}{2}}} > \Psi^{\dagger}(y_{j}) \|\frac{y_{j}}{\|y_{j}\|} = \frac{2\Psi(y_{j})}{\|y_{j}\|}$$

is bounded away from 0. It follows that $(\lambda(y_j))$ is a bounded sequence and therefore by (ii) of Proposition 2.10, $S(y_j)$ has a convergent subsequence.

To get the compactness of K, we use a variant of an argument of BENCI [15]. Let $B \subset X$ be bounded. Without loss of generality, $B = B_R$, a ball of radius R about 0. By 3° of Proposition 3.1, $n([0,1] \times B_R) \subset B_{R+1}$. Therefore $S(n([0,1] \times B_R)) \subset \overline{S(B_{R+1})}$ which is compact. If

$$Y \equiv \{aw \mid a \in \{0,1\}, w \in \overline{S(B_{R+1})}\}$$
,

then Y is compact as is \hat{Y} , its closed convex hull. Recalling that $\omega(\xi) \in [0,1]$, it follows that for each $\tau \in [0,s]$ and $z \in B_p$,

$$Z \equiv \{\exp(\int_{S}^{\tau} \omega(n(r,z))dr)\}S(n(\tau,z)) \in Y$$
.

Hence for $s \in [0,1]$,

$$\int_0^B z d\tau \in \hat{Y},$$

K is compact, and the proof is complete.

Now define $M^+ \equiv X^+ \cap M$,

$$w = x^0 \cdot x^- \cdot span\{\varphi\}$$

where $\varphi = (\sin t)e_1 - (\cos t)e_{n+1} \in X^+$. Set $M^- = W \cap M$. Define

(3.6)
$$\underline{\alpha} \equiv \inf_{z \in M^+} A(z)$$

and

(3.7)
$$\vec{a} \equiv \sup_{z \in M} A(z) .$$

Proposition 3.8: $0 < \alpha < \overline{\alpha} < \infty$.

Proof: If $z \in X^+$, $A(z) = \|z\|^2$. If $\underline{\alpha} = 0$, $0 \in \overline{M}^+$, $= M^+$. Since by (ii) of Proposition 2.18, M bounds a neighborhood of 0 in X, this is impossible and $\underline{\alpha} > 0$. Next observe that

$$\underline{\alpha} \le \inf_{\text{span}\{w\} \cap M} A(z) \le \sup_{M} A(z) = \overline{\alpha}.$$

Finally note that if $z \in M^-$, $z = r(z)\phi + z^0 + z^-$. Therefore since $A(\phi) = \pi$,

(3.9)
$$A(z) \le r^2(z)\pi$$
.

Since $z \in M$, by (2.2),

(3.10)
$$\frac{1}{2\pi} \int_{0}^{2\pi} m|z|^{2} dt \leq 1 = \Psi(z) .$$

Hence (3.10) and the orthogonality of X^{O} , X^{\pm} in L^{2} imply

(3.11)
$$\frac{2\pi}{m} > r(z)^2 \int_0^{2\pi} |\phi|^2 dt = 2\pi r(z)^2.$$

Therefore (3.7) and (3.11) show

$$(3.12) \qquad \qquad \bar{\alpha} \leq \pi r(z)^2 \leq \frac{\pi}{z}.$$

Our goal is to obtain a critical value c of $A|_{M}$ via a minimax argument. The class of maps which will be used to define c can now be introduced. Let Γ denote the set of $h \in C(X,X)$ satisfying the following three conditions:

$$(\Gamma_1)$$
 h(z) = z if A(z) $\not\in [0,\overline{\alpha}+1]$ or if $|\Psi(z)-1| > \frac{1}{2}$.

 (Γ_2) $P^-h(z) = e^{\phi(z)}z^- + Q(z)$ where $\phi \in C(X, \mathbb{R}^+)$, $0 < \phi < \gamma$, γ depending on h, and Q is compact.

$$(\Gamma_3)$$
 h: M+M.

Remark 3.13: Observe that Γ is closed under composition. Moreover $1^{O}-5^{O}$ of Proposition 3.1 imply $\eta(1, \cdot) \in \Gamma$ provided that $0 \le c - \overline{\epsilon} \le c + \overline{\epsilon} \le \overline{\alpha} + 1$. This inequality holds in particular if $c \in (0, \overline{\alpha}]$ and $\overline{\epsilon}$ is chosen appropriately.

The mappings in I satisfy an important intersection property.

Proposition 3.14: If $h \in \Gamma$, then $h(M^-) \cap M^+ \neq \emptyset$.

This proposition will be proved in §4. Assuming it for now, define

(3.15)
$$c = \inf_{h \in \Gamma} \sup_{z \in M^-} A(h(z)).$$

Proposition 3.16: $\alpha \le c \le \overline{\alpha}$ and c is a critical value of $A|_{M}$.

Proof: If $h \in \Gamma$, by Proposition 3.14,

(3.17)
$$\sup_{z \in M^-} A(h(z)) > \sup_{\zeta \in h(M^-) \cap M^+} A(\zeta) > \inf_{\zeta \in M^+} A(\zeta) = \underline{\alpha}.$$

Since (3.17) holds for all $h \in \Gamma$, $c > \underline{\alpha}$. On the other hand $h(z) \equiv z \in \Gamma$. Consequently (3.18) $c \leq \sup_{z \in M} A(z) = \overline{\alpha}.$

To prove that c is a critical value of $\mathbb{A}|_{M'}$ suppose on the contrary that $\mathbb{K}_{\mathbb{C}} = \beta$. Then if $\overline{\epsilon} < \min(\frac{1}{2}\overline{\alpha}, 1)$, Remark 3.13 shows $n(1, \cdot)$ as determined from Proposition 3.1 belongs to Γ . Choose $h \in \Gamma$ such that

(3.19)
$$\sup_{z \in \mathbb{N}^{-1}} \lambda(h(z)) \le c + \varepsilon$$

where ϵ is obtained from Proposition 3.1. Since $\eta(1,h) \in \Gamma$, by 6° of Proposition 3.1, (3.15), and (3.19),

(3.20)
$$c \leq \sup_{z \in M} \lambda(\eta(1,h(z))) \leq c - \varepsilon$$

a contradiction and the proposition is proved.

Completion of proof of Theorem 1.2 for $H \in \mathbb{C}^2$: By Proposition 3.16, there is a $z \in M$ such that A(z) = c and

$$(3.21) \qquad (\lambda^{\dagger}(z) - \lambda(z)\Psi^{\dagger}(z))(\zeta) = 0$$

for all $\zeta \in X$. This implies z is a classical solution of (2.4). Indeed $z \in X$ implies $\widetilde{H}_Z(z) \in L^2(S^1, \mathbb{R}^{2n})$. Moreover by (H_3) , $\widetilde{H}_p(p(t), q(t))$ is odd about 0 and π and $\widetilde{H}_q(p(t), q(t))$ is even about 0 and π . Taking $\zeta \equiv 1$ in (3.21) yields $[\widetilde{H}_Z(z(t))] = 0$ where

$$[w] = \frac{1}{2\pi} \int_{0}^{2\pi} w(t) dt$$
.

Since $[\hat{H}_{p}(p(t),q(t))] = 0$, Fourier expansion shows the equations

(3.22) (i)
$$\frac{d\mathbf{p}}{dt} = -\lambda(z) \tilde{\mathbf{H}}_{\mathbf{q}}(\mathbf{p}(t), \mathbf{q}(t))$$
(ii)
$$\frac{d\mathbf{Q}}{dt} = \lambda(z) \tilde{\mathbf{H}}_{\mathbf{p}}(\mathbf{p}(t), \mathbf{q}(t))$$

have a unique solution $Z=(P,Q)\in X\cap W^{1,2}(S^1,R^{2n})$ with [Q]=[q]. For smooth $\zeta=(\varphi,\psi)\in X$, taking the inner product of (3.22) (i) with φ and (ii) with ψ gives

$$(3.23) \qquad (\lambda'(z) - \lambda(z) \forall'(z))(\zeta) = 0.$$

Comparing (3.21) and (3.23) shows

(3.24)
$$\int_{0}^{2\pi} \{(p-P) \cdot \sqrt[4]{q} - (q-Q) \cdot \sqrt[4]{q}\} dt = 0$$

for all smooth $(\phi,\psi)\in X$ where [p-P]=0=[q-Q]. Hence $z=z\in W^{1/2}(S^1,\mathbb{R}^{2n})\subset C(S^1,\mathbb{R}^{2n}).$ Thus (3.22) shows $z\in C^1(S^1,\mathbb{R}^{2n})$ and is a classical solution of (2.4). But (2.4) is a Hamiltonian system so $\widetilde{H}(z(t))$ is independent of t. Consequently

(3.25)
$$1 = \Psi(z) = \frac{1}{2\pi} \int_{0}^{2\pi} \overline{H}(z(t)) dt = \overline{H}(z(t))$$

so $z(t) \in \overline{H}^{-1} = H^{-1}(1)$. Finally Lemma 2.3 gives a solution of the desired type of (HS).

§4. The Intersection Theorem and the General Case of Theorem 1.2

In this section we will prove the intersection theorem: Proposition 3.14, and obtain the C^1 case of Theorem 1.2. For the former result, the following technical result is required:

Proposition 4.1: Let V be a k dimensional subspace of \mathbb{R}^n and \mathbb{V}^1 its orthogonal complement. Suppose h $\in C(\mathbb{R}^n,\mathbb{R}^n)$ and satisfies

 (h_4) h = id on v^1

and

(h₂) there is an R > 0 such that h = id on $\mathbb{R}^{n} \backslash \mathbb{B}_{R}$. Let $\psi \in \mathbb{C}^{1}(\mathbb{R}^{n}, \mathbb{R})$ and satisfy

 (ψ_1) $\psi(0) = 0$ and $x \cdot \psi^*(x) > 0$ for $x \neq 0$

and

 (ψ_2) there is a $\rho \in (0,R)$ such that $\psi^{-1}(\rho) \subseteq B_R$.

Let $v \in V \cap \partial B_1$ and set $Y = \text{span}\{v\} \oplus V^{\perp}$. Then there is a $\xi \in Y$ such that

$$\psi(\xi) = \rho \text{ and } h(\xi) \in V.$$

Proof: Let $Q \equiv \{rv \mid 0 < r < R\} \in (B_R \cap V^1)$ so $Q \subset Y$. We will find $\xi \in Q$. Let P and P^1 denote the orthogonal projectors of R^n onto V and V^1 . Solving (4.2) for $\xi \in Y$ is equivalent to finding $\xi \in Y$ such that

(4.3) (i) $\psi(\xi) = \rho$

(ii) $P^{\perp}h(\xi) = 0$.

For $y \in Y$, set

$$\Phi(y) = (\psi(y), P^{\perp}h(y)) .$$

Identifying $\mathbb{R} \times \mathbb{Y}$ with $\mathbb{R} \times \mathbb{R}^{n-k}$ and \mathbb{Q} with a subset thereof, Φ can be considered to be a continuous map of $\mathbb{R} \times \mathbb{R}^{n-k}$ into itself. Any zero of Φ is a solution of (4.3). Consider $d(\Phi,\mathbb{Q},(\rho,0))$, the Brouwer degree of Φ with respect to the bounded open set \mathbb{Q} and the point $(\rho,0)$. We will show this degree equals 1 and therefore (4.3) has a solution in \mathbb{Q} . In order for the degree to be defined, it is necessary that $\Phi \neq (\rho,0)$ on $\partial \mathbb{Q}$. Writing $\mathbb{Y} \in \mathbb{Q}$ as $(r,w) \in \mathbb{R} \times \mathbb{R}^{n-k}$, if r=0, by (h_1) , h=id

so by (ψ_1) , $\Phi(y) = (\psi(w), w) \neq (\rho, 0)$. If r = R, (h_2) implies h = id and $\Phi(y) = (\psi(Rv + w), w) \neq (\rho, 0)$ since $\psi(Rv) > \rho$ via $(\psi_1) - (\psi_2)$. Finally if |w| = R, (h_2) implies h = id and $\Phi(y) = (\psi(rv + w), w) \neq (\rho, 0)$. Therefore $\Phi(\Phi, Q, (\rho, 0))$ is defined. We claim

$$d(\Phi,Q(\rho,0)) = d(id,Q,(\rho,0)).$$

Since $(\rho,0) \in Q$ via (ψ_2) ,

Season serential

$$d(id,Q,(\rho,0)) = 1$$

and the proof is complete. To verify (4.4), consider the homotopy

$$\Phi_{\theta}(y) = (\theta r + (1 - \theta)\psi(y), p^{\perp}h(y))$$

for $y \in \overline{Q}$ and $\theta \in [0,1]$. Arguing as for θ , if r = 0, $\theta_{\theta}(y) = ((1-\theta)\psi(w), w) \neq (\rho, 0); \text{ if } r = R, \ \theta_{\theta}(y) = (\theta R + (1-\theta)\psi(RV + w), w) \neq (\rho, 0)$ since R, $\psi(RV) > \rho_{1}$ if |w| = R, $\theta_{\theta}(y) = (\theta r + (1-\theta)\psi(rV + w), w) \neq (\rho, 0)$. Since $\theta_{0}(y) = \theta(y) \text{ and } \theta_{1}(y) = y \text{ on } \theta_{Q}, \text{ the homotopy invariance property of Brouwer degree yields (4.4) and the proof is complete.}$

Remarks 4.5: (i) The same hypotheses and an appropriately modified Φ also yield an $y \in Y$ such that $h(y) \in V \cap \psi^{-1}(\rho)$. This fact can be used to obtain a critical value of $\lambda|_{M}$ as a maximin rather than a minimax. (ii) An examination of the proof of Proposition 4.1 shows we need merely take $\psi \in C(\mathbb{R}^n, \mathbb{R})$ and weaken (ψ_1) to $\psi(0) = 0$ and $\psi(\mathbb{R}^n) > \rho$ for some $v \in V \cap \partial B_1$. Also at the expense of redefining B_R , V and V^L can be any complementary subspaces of \mathbb{R}^n .

Now we can give the:

Proof of Proposition 3.14: Let X_1^{\pm} denote the subspaces of X^{\pm} defined by restricting j to 1 < j < i in the definition of X^{\pm} . Let $X_1 = X^0 \oplus X_1^{\pm} \oplus X_1^{-}$ and let P_1 denote the projector of X onto X_1 . Define $h_1(z) \equiv P_1h(z)$. By (2.2),

$$\Psi(z) > \frac{m}{2\pi} \|z\|_{L^2}^2.$$

Hence $\Psi(z) + \infty$ as $\|z\| + \infty$ uniformly for $z \in X_1$. In particular there is an $R_1 > 0$ such that $\Psi(z) > 2$ if $z \in X_1$ and $\|z\| > R_1$. Therefore if $h \in \Gamma$, by (Γ_1) ,

 $\begin{aligned} h(z) &= z = h_{\underline{i}}(z) & \text{ if } z \in X_{\underline{i}} & \text{ and } \exists z \ni R_{\underline{i}}. & \text{Moreover if } z \in X^0 \oplus X_{\underline{i}}^-, A(z) \le 0. \end{aligned}$ Therefore by (Γ_1) again, $h_{\underline{i}}(z) = z$ on $X^0 \oplus X_{\underline{i}}^-. \text{ Since } \Psi|_{X_{\underline{i}}} & \text{satisfies } (\psi_1) - (\psi_2) \text{ of } Proposition 4.1, this result implies there is a <math>z_{\underline{i}} \in W_{\underline{i}} \cap M^-$ such that $h_{\underline{i}}(z_{\underline{i}}) \in X_{\underline{i}}^+$ where $W_{\underline{i}} \equiv X^0 \oplus X_{\underline{i}}^- \oplus \text{span}\{\varphi\}.$

We claim (z_1) is bounded in X. Let $z_1=z_1^0+z_1^-+z_1^+$. As was noted earlier, these components of z_1 are mutually orthogonal in L^2 . By Proposition 2.18 (iii), (z_1) is bounded in $L^2(S^1,\mathbb{R}^{2n})$. Since X^0 is finite dimensional and z_1^+ is a bounded multiple of φ , $(z_1^0+z_1^+)$ is bounded in X. If (z_1^-) is unbounded, $A(z_1)=\|z_1^+\|^2-\|z_1^-\|^2+-\infty$. Therefore for large i, $h(z_1)=z_1=h_1(z_1)\in W_1\cap M^-\cap X_1^+$, i.e. $z_1=z_1^+$ for large i and (z_1) is bounded in X. It is clear that $(z_1^0+z_1^+)$ possesses a convergent subsequence. By (Γ_2) ,

$$0 = P^{-}h_{i}(z_{i}) = e^{\varphi(z_{i})}z_{i}^{-} + P^{-}P_{i}Q(z_{i})$$

OF

(4.6)
$$z_{1}^{-} = -e^{-\varphi(z_{1})} p^{-} p_{1} Q(z_{1})$$
.

Since Q is compact and $0 \le \varphi(\cdot) \le \gamma$, (4.6) shows (z_1^-) also has a convergent subsequence along which $z_1^- + z \in M^-$. Since h is continuous, $P_1h(z_1^-) + h(z) \in X^+$. Finally by (Γ_3^-) , $h(z) \in M^+$. Thus $h(M^-) \cap M^+ \neq \emptyset$ and the proof is complete.

Next we will give the

<u>Proof of Theorem 1.2 for $H \in C^1(\mathbb{R}^{2n},\mathbb{R})$ </u>: Let $(\vec{\mathbb{H}}_k)$ be a sequence of C^2 functions which are homogeneous of degree 2, satisfy (\mathbb{H}_3) , and converge to $\vec{\mathbb{H}}$ in C^1 uniformly in a neighborhood of S^{2n-1} . The C^2 version of Theorem 1.2 implies there is a $z_k \in X$ which is a classical solution of

(4.7)
$$\dot{z}_{k} = \lambda_{k} \hat{H}_{kz}(z_{k})$$

and $z_k(0) \in \overline{H}_k^{-1}(1)$. Equation (4.7) implies that

(4.8)
$$c_{k} \equiv A(z_{k}) = \frac{\lambda_{k}}{2} \int_{0}^{2\pi} z_{k} \cdot \tilde{R}_{kz}(z_{k}) dt = 2\pi \lambda_{k}.$$

By Proposition 3.16, $c_k > 0$. Suppose that (c_k) is bounded away from 0 and -. Then (4.8) shows the same is true for (λ_k) . Therefore (4.7) provides L^{∞} bounds for \dot{z}_k and (4.7) and the Arzela-Ascoli Theorem imply a subsequence of (λ_k, z_k) converge in $\mathbb{R}^+ \times \mathbb{C}^1$ to a solution (λ, ζ) of (2.4). Following the \mathbb{C}^1 version of the proof of Lemma 2.3 then gives a solution of (HS) of the desired type. (Now β in (2.6) is merely continuous so (2.6) need not have a unique solution but any solution will suffice.)

It remains to get the bounds for c_k . By Proposition 3.16, there are constants $\underline{\alpha}_k$, $\bar{\alpha}_k$ such that

$$(4.9) \underline{\alpha}_{\nu} < c_{\nu} < \overline{\alpha}_{\nu}$$

where $\underline{\alpha}_k$, $\overline{\alpha}_k$ are defined in (3.6), (3.7) with $M=M_k$. By (3.12), $\overline{\alpha}_k \leq \pi m_k^{-1}$ where m_k is defined in (2.1). Since $\overline{H}_k + \overline{H}$ uniformly on S^{2n-1} , $m_k + m$ so for large k, $m_k > \pi/2$ and

$$c_{k} \leq \overline{\gamma}_{k} \leq 2\pi m^{-1}.$$

Thus (c_k) is bounded away from ∞ . To get a lower bound for c_k , recall that by (2.2), $\widetilde{\mathrm{H}}_k(z) \leq \mathrm{M}_k |z|^2$

with M_k defined in (2.1) and $M_k \to M$ as $h \to \infty$. Hence $M_k \le 2M$ for large k and if $z \in M_k$,

$$1 = \frac{1}{2\pi} \int_{0}^{2\pi} \tilde{H}_{k}(z) dt \leq \frac{M}{\pi} \int_{0}^{2\pi} |z|^{2} \equiv \Psi(z) .$$

If $z \in \P^{-1}(1)$, there is an $a_k(z) > 1$ such that

$$\Psi_{k}(a_{k}(z)z) = \frac{1}{2\pi} \int_{0}^{2\pi} \vec{H}_{k}(a_{k}(z)z)dt = 1$$
.

Therefore

(4.11)
$$\underline{\alpha}_{+} = \inf_{z \in \Psi^{-1}(1) \cap K^{+}} A(z) \leq \underline{\alpha}_{k}$$

for large k. The argument of Proposition 3.8 shows $\alpha_*>0$. Hence (4.9) and (4.11) show (c_k) is bounded from below and the proof is complete.

REFERENCES

- [1] SEIFERT, H., Periodische Bewegungen mechanischen Systeme, Math. Z., <u>51</u> (1948), 197-216.
- [2] RUIZ, O. R., Existence of brake orbits in Finsler mechanical systems, Springer Lec. Notes in Math., <u>597</u> (1977), 542-567.
- [3] WEINSTEIN, A., Periodic orbits for convex Hamiltonian systems, Ann. Math., 108 (1978), 507-518.
- [4] RABINOWITZ, P. H., Periodic solutions of Hamiltonian systems, Comm. Pure Appl. Math., 31, (1978), 157-184.
- [5] RABINOWTTZ, P. H., Periodic solutions of a Hamiltonian system on a prescribed energy surface, J. Diff. Eq., 33 (1979), 336-352.
- [6] CLARKE, F., A classical variational principle for periodic Hamiltonian trajectories, Proc. Am. Math. Soc., 76 (1979), 186-188.
- [7] GLUCK, H. and W. ZILLER, Existence of periodic solutions of conservative systems,

 Seminar on Minimal Submanifolds, Princeton University Press, 1983, 65-98.
- [8] HAYASHI, K., Periodic solution of classical Hamiltonian systems, Tokyo J. Math., 6
 (1983), 473-486.
- [9] BENCI, V., Closed geodesics for the Jacobi metric and periodic solutions of prescribed energy of natural Hamiltonian systems, to appear, Ann. Inst. H. Poincare, Analyse nonlineaire.
- [10] van GROESEN, E. W. C., Existence of multiple normal mode trajectories on convex energy surfaces of even classical Hamiltonian systems, to appear, J. Diff. Eq.
- [11] GIRARDI, M., Multiple orbits for Hamiltonian systems on starshaped surfaces with symmetries, Ann. Inst. H. Poincare, Analyse nonlineaire, 1 (1984), 285-294.
- [12] EKELAND, I. and J. M. LASRY, On the number of periodic trajectories for a Hamiltonian flow on a convex energy surface, Ann. Math., 112 (1980), 283-319.

- [13] RABINOWITZ, P. H., A variational method for finding periodic solutions of differential equations, Nonlinear Evolution Equations (M. G. CRANDALL, editor), Academic Press, New York (1978), 225-251.
- [14] RABINOWITZ, P. H., Variational methods for nonlinear eigenvalue problems, Proc. Sym. Eigenvalues of Nonlinear Problems, (G. PRODI, editor), Edizioni Cremonese, Rome (1974), 139-195.
- [15] BENCI, V., On critical point theory for indefinite functionals in the presence of symmetries, Trans. Amer. Math. Soc., <u>274</u> (1982), 533-572.

PHR: scr

REPORT DOCUMEN	READ INSTRUCTIONS BEFORE COMPLETING FORM		
. REPORT NUMBER	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER	
2901	ł		
4. TITLE (and Subtitle)		5. Type of Report a Period Covered Summary Report - no specific	
ON THE EXISTENCE OF PERIODIC SOLUTIONS FOR A		reporting period	
CLASS OF SYMMETRIC HAMILTON	6. PERFORMING ORG. REPORT NUMBER		
7. AUTHOR(e)	8. CONTRACT OR GRANT NUMBER(#)		
Paul H. Rabinowitz		DAAG29-80-C-0041 MCS-8110556	
PERFORMING ORGANIZATION NAME AND	10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS		
Mathematics Research Center	Work Unit Number 1 -		
610 Walnut Street	Wisconsin	Applied Analysis	
Madison, Wisconsin 53705			
1. CONTROLLING OFFICE NAME AND ADD	12. REPORT DATE		
See Item 18 below.		January 1986	
		13. NUMBER OF PAGES	
		20	
14. MONITORING AGENCY NAME & ADDRES	18. SECURITY CLASS. (of this report)		
		UNCLASSIFIED	
		15a. DECLASSIFICATION/DOWNGRADING	
6. DISTRIBUTION STATEMENT (of this Rep	ort)	<u> </u>	
Approved for public release;	distribution unlimited.		
· · · · · · · · · · · · · · · · · ·			

- 17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)
- 18. SUPPLEMENTARY NOTES
- U. S. Army Research Office
- P. O. Box 12211

Research Triangle Park

North Carolina 27709

19. KEY WORDS (Continue on reverse side if necessary and identify by block number)

periodic solution

Hamiltonian system

minimax methods

20. ABSTRACT (Continue on reverse side if necessary and identify by block number)

The main result in this paper is:

Theorem: If $H \in C^1(\mathbb{R}^{2n}, \mathbb{R})$ and satisfies

 (H_1) $H^{-1}(1)$ bounds a starshaped neighborhood of 0 in \mathbb{R}^{2n} ,

 (H_2) $z \cdot H_z \neq 0$ for all $z \in H^{-1}(1)$,

DD FORM 1473 EDITION OF 1 NOV 65 IS OBSOLETE

UNCLASSIFIED

National Science Foundation

Washington, DC 20550

20. ABSTRACT - cont'd.

 (H_3) H(p,q) = H(-p,q) for all $p,q \in \mathbb{R}^n$, then there is a T>0 such that the Hamiltonian system

(HS)
$$\dot{z} = J_{H_z}(z), \quad J = \begin{pmatrix} 0 & -id \\ id & 0 \end{pmatrix}$$

possesses a T periodic solution $(p(t),q(t)) \in H^{-1}(1)$ with p odd about 0 and T/2 and q even about 0 and T/2.

The proof involves a new existence mechanism which should be useful in other situations.

6-86